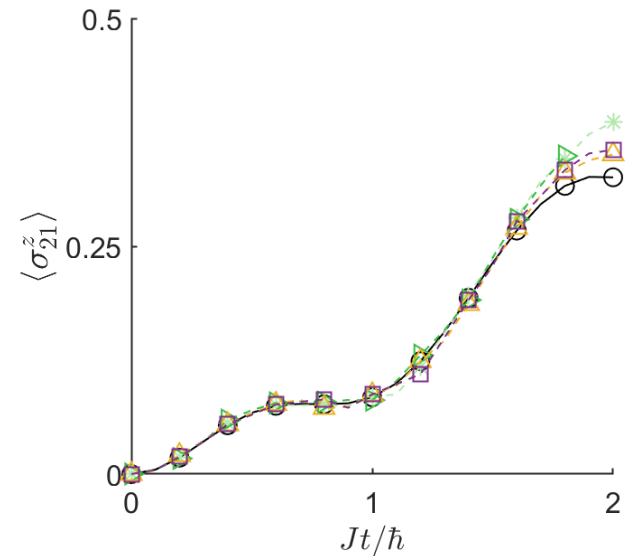
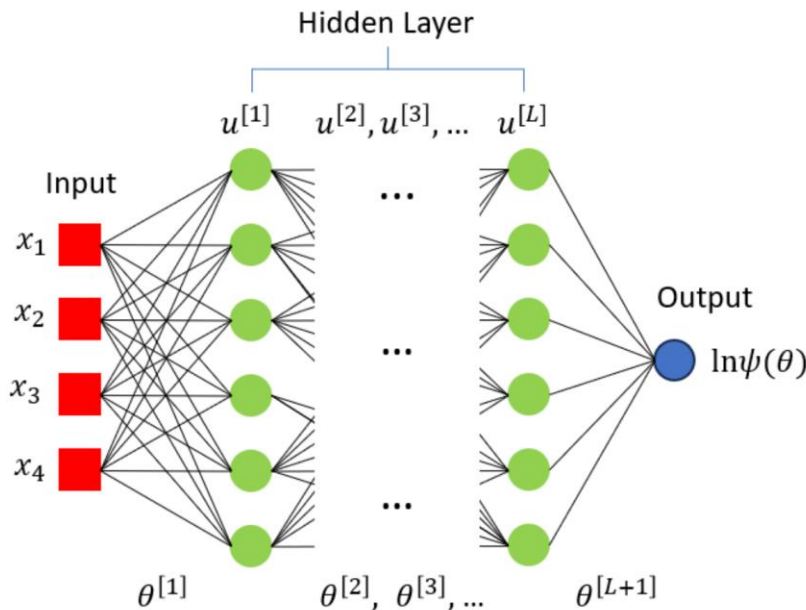


Paths towards time evolution with larger neural-network quantum states



Machine learning for quantum technology
Erlangen 2024

Dario Poletti

Singapore University of Technology and Design
Centre for Quantum Technologies
Majulab



Zhang et al., arxiv:2406.03381

Paths towards time evolution with larger neural-network quantum states

Principal Investigator
Dario Poletti

PostDocs
Noufal Jaseem
Wang Dingzu
Zhang Wenxuan

PhD Students
Rebecca Erbanni
Gauthameswar
Apimuk Sornsaeng

Collaborators on
these topics

Guo Chu
Itai Arad
Lu Wei
Kavan Modi

Remmy Zen
Stephane Bressan
Christian Miniatura
Xu Xiansong
Xing Bo



Xu Xiansong



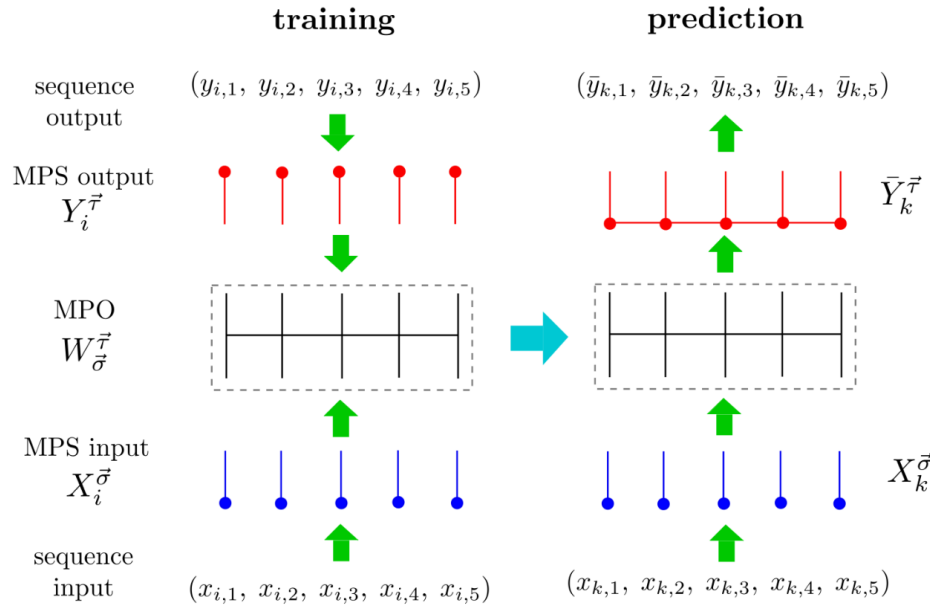
Xing Bo



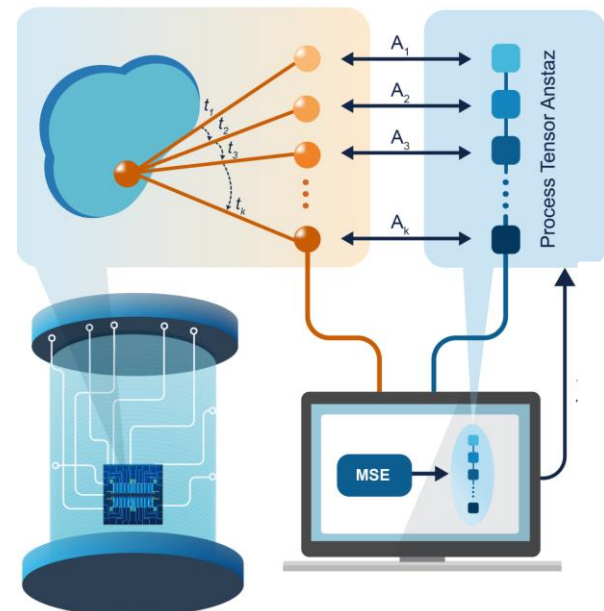
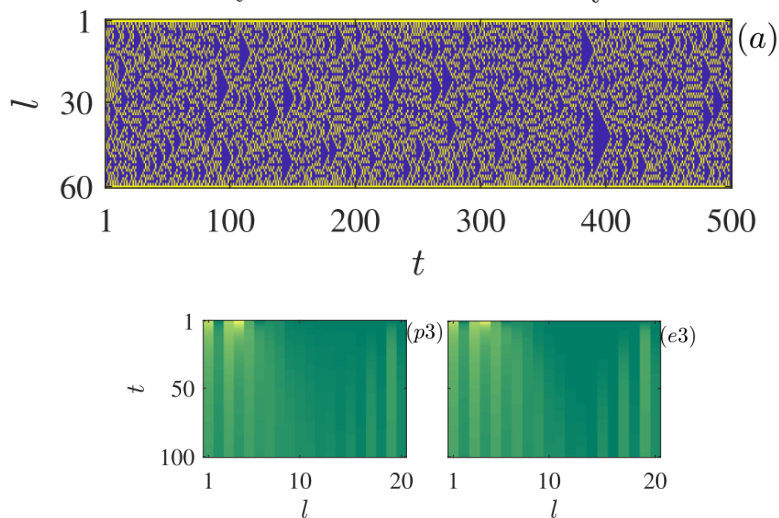
Zhang Wenxuan



Using MPOs for sequence-to-sequence learning



Guo, et al., PRE 2018
 C. Guo et al., PRA 2020
 H.P. Casagrande et al., arXiv:2404.11768
 X. Zhang, et al., arXiv:2312.06062
 C. Guo, et al., manuscript in preparation

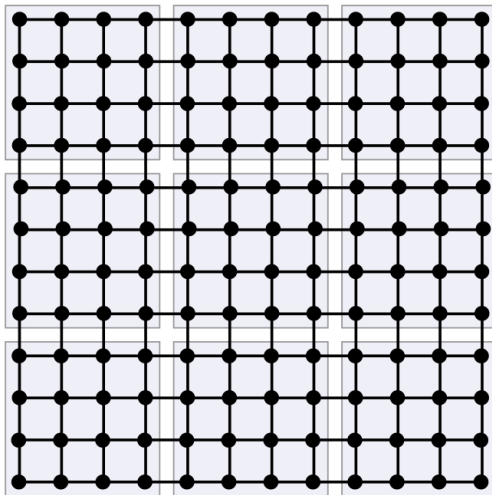


Block belief propagation and tensor networks

Merging **tensor networks**
with **belief propagation**

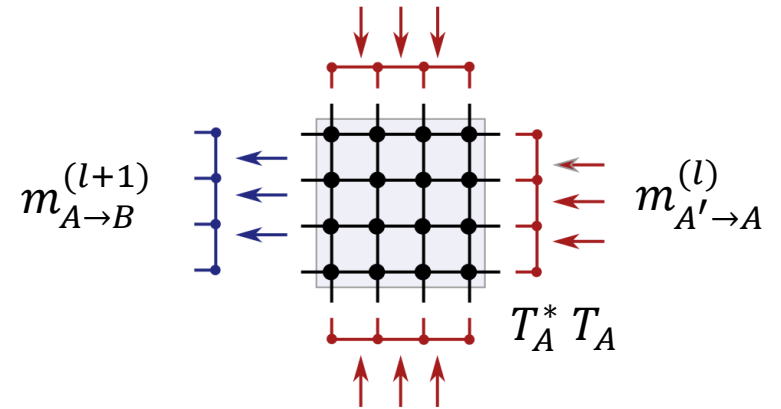
Divide the system into blocks

(a)



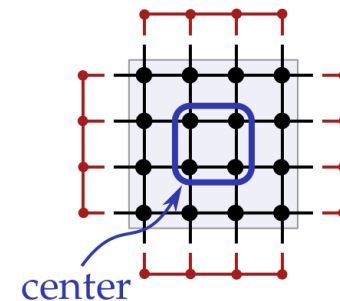
Update the message coming from one block

(b) blockBP iteration



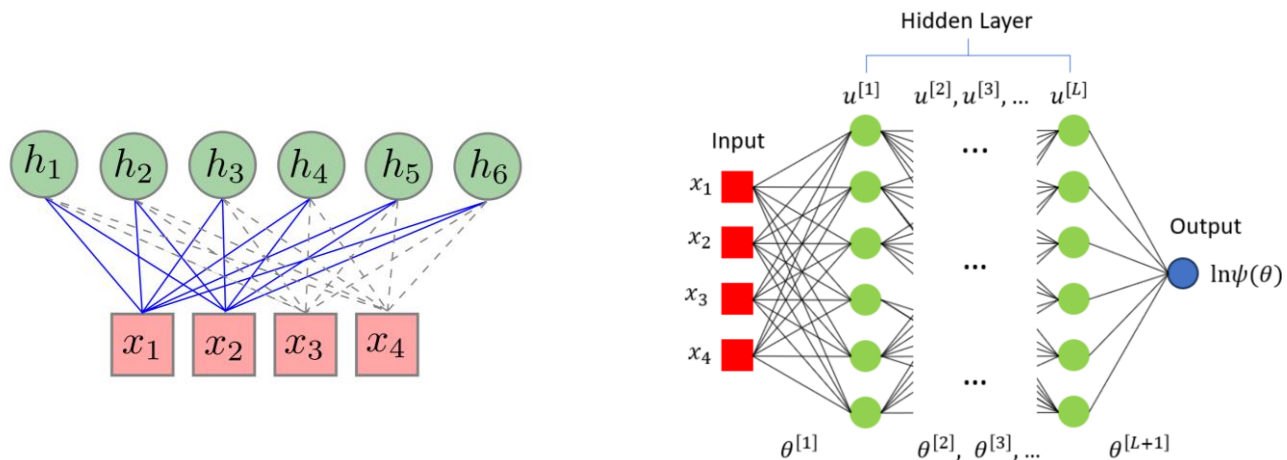
Compute local quantities using messages from other blocks

(c) Local environment



Plan of the presentation

- 1) Time evolution from overlap
 - Loss function
 - Optimization approaches
- 2) Holomorphic neural networks
- 3) Quantum quench in the tilted Ising model
- 4) Conclusions and outlook



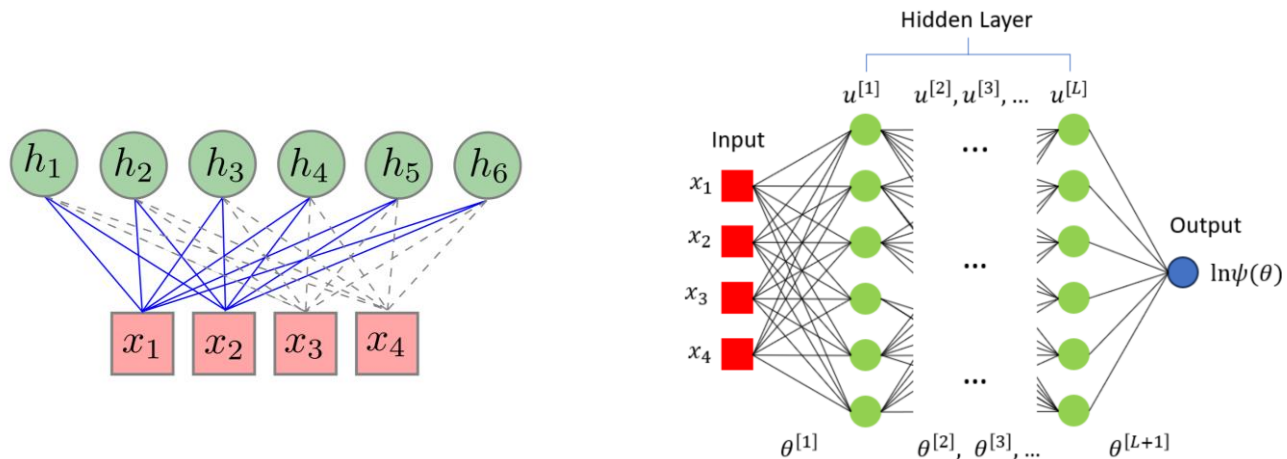
Plan of the presentation

- 1) Time evolution from overlap
 - Loss function
 - Optimization approaches

2) Holomorphic neural networks

3) Quantum quench in the tilted Ising model

4) Conclusions and outlook



Time evolution with NQS

From Carleo and Troyer's 2017 paper we have been presented with one approach to simulate time evolution of many-body quantum systems with tensor networks

$$\frac{d\boldsymbol{\theta}(t)}{dt} = -i \mathbf{S}^{-1}(t) \mathbf{F}(t)$$

where θ are the network parameters, and the equation is then derived from Schrodinger's equation. Also known as tVMC.

Here $\mathbf{S} = \mathbf{X}\mathbf{X}^\dagger$ and $\mathbf{F} = \mathbf{X}\mathbf{f}$ with

$$X_{m,k} = \frac{1}{\sqrt{N_s}} \left(\left(\frac{\partial \ln[\psi(\mathbf{x}_k; \boldsymbol{\theta})]}{\partial \theta_m} \right)^* - \left\langle \frac{\partial \ln[\psi(\mathbf{x}; \boldsymbol{\theta})]}{\partial \theta_m} \right\rangle^* \right)$$

and

$$f_k = \frac{1}{\sqrt{N_s}} (E_{loc}(\mathbf{x}_k) - \langle E_{loc}(\mathbf{x}) \rangle)$$

Time evolution with NQS

A different approach aims to find directly the neural network which best approximates the evolved one. Also known as p-tVMC.

For instance,

$$\|\psi[\theta_{n+1}] - \Phi^{\Delta t}(\psi[\theta_n])\|$$

Time evolution with NQS

A different approach aims to find directly the neural network which best approximates the evolved one. Also known as p-tVMC.

For instance,

$$\|\psi[\theta_{n+1}] - \Phi^{\Delta t}(\psi[\theta_n])\|$$

This can be done in various ways, either expanding the unitary evolution

$$\psi[\theta_{n+1}] \approx \psi[\theta_n] - i\Delta t H \left(\frac{\psi[\theta_{n+1}] + \psi[\theta_n]}{2} \right)$$

Time evolution with NQS

A different approach aims to find directly the neural network which best approximates the evolved one. Also known as p-tVMC.

For instance,

$$\|\psi[\theta_{n+1}] - \Phi^{\Delta t}(\psi[\theta_n])\|$$

This can be done in various ways, either expanding the unitary evolution

$$\psi[\theta_{n+1}] \approx \psi[\theta_n] - i\Delta t H \left(\frac{\psi[\theta_{n+1}] + \psi[\theta_n]}{2} \right)$$

or not

$$\mathcal{I}_{\text{loc}}(\sigma, \eta) = 1 - \frac{\langle \sigma | \mathcal{U} | \Psi_{\theta} \rangle \langle \eta | \mathcal{U}^{\dagger} | \Psi_{\bar{\theta}} \rangle}{\langle \sigma | \Psi_{\bar{\theta}} \rangle \langle \eta | \Psi_{\theta} \rangle}$$

Neural network quantum states: **time evolution**

Cost function

$$C_{\psi_t, \psi_{t'}}^U = \frac{\langle \psi_{t'} | U | \psi_t \rangle \langle \psi_t | U^\dagger | \psi_{t'} \rangle}{\langle \psi_{t'} | \psi_{t'} \rangle \langle \psi_t | \psi_t \rangle}$$

Neural network quantum states: **time evolution**

Cost function

$$C_{\psi_t, \psi_{t'}}^U = \frac{\langle \psi_{t'} | U | \psi_t \rangle \langle \psi_t | U^\dagger | \psi_{t'} \rangle}{\langle \psi_{t'} | \psi_{t'} \rangle \langle \psi_t | \psi_t \rangle}$$

Forward evolution

Backward evolution

Neural network quantum states: **time evolution**

Cost function

$$C_{\psi_t, \psi_{t'}}^U = \frac{\langle \psi_{t'} | U | \psi_t \rangle \langle \psi_t | U^\dagger | \psi_{t'} \rangle}{\langle \psi_{t'} | \psi_{t'} \rangle \langle \psi_t | \psi_t \rangle}$$

An effective Hamiltonian



Neural network quantum states: **time evolution**

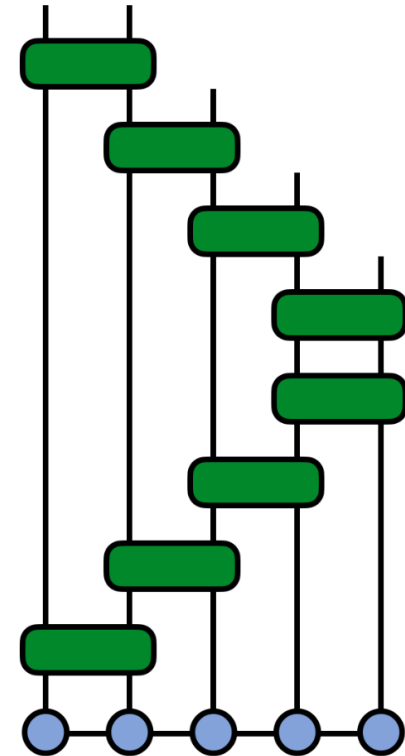
Cost function

$$C_{\psi_t, \psi_{t'}}^U = \frac{\langle \psi_{t'} | U | \psi_t \rangle \langle \psi_t | U^\dagger | \psi_{t'} \rangle}{\langle \psi_{t'} | \psi_{t'} \rangle \langle \psi_t | \psi_t \rangle}$$

An effective Hamiltonian

$$U(dt) = \prod_{i=1}^{L-1} U_{2,i}(dt/2) \prod_{j=L-1}^1 U_{2,j}(dt/2)$$

Suzuki-Trotter decomposition



In this way the U used at each step is sparse

Neural network quantum states: approaches to train larger networks

Cost function

$$\begin{aligned} C_{\psi_t, \psi_{t'}}^U &= \sum_{\mathbf{x}, \mathbf{y}} P_{\psi_{t'}}(\mathbf{x}) P_{\psi_t}(\mathbf{y}) E_{\psi_t \psi_{t'}}^U(\mathbf{x}) E_{\psi_{t'} \psi_t}^U(\mathbf{y}) \\ &= \sum_{\mathbf{x}} P_{\psi_{t'}}(\mathbf{x}) E_{loc}^U(\mathbf{x}), \end{aligned}$$

this is readily amenable to use Stochastic Reconfiguration for the time evolution.

$$\boldsymbol{\theta} = \boldsymbol{\theta} - \gamma \mathbf{S}^{-1} \mathbf{F}$$

Neural network quantum states: approaches to train larger networks

Cost function

$$\begin{aligned} C_{\psi_t, \psi_{t'}}^U &= \sum_{\mathbf{x}, \mathbf{y}} P_{\psi_{t'}}(\mathbf{x}) P_{\psi_t}(\mathbf{y}) E_{\psi_t \psi_{t'}}^U(\mathbf{x}) E_{\psi_{t'} \psi_t}^U(\mathbf{y}) \\ &= \sum_{\mathbf{x}} P_{\psi_{t'}}(\mathbf{x}) E_{loc}^U(\mathbf{x}), \end{aligned}$$

this is readily amenable to use Stochastic Reconfiguration for the time evolution.

$$\boldsymbol{\theta} = \boldsymbol{\theta} - \gamma \mathbf{S}^{-1} \mathbf{F}$$

But inverting the S matrix becomes more and more demanding

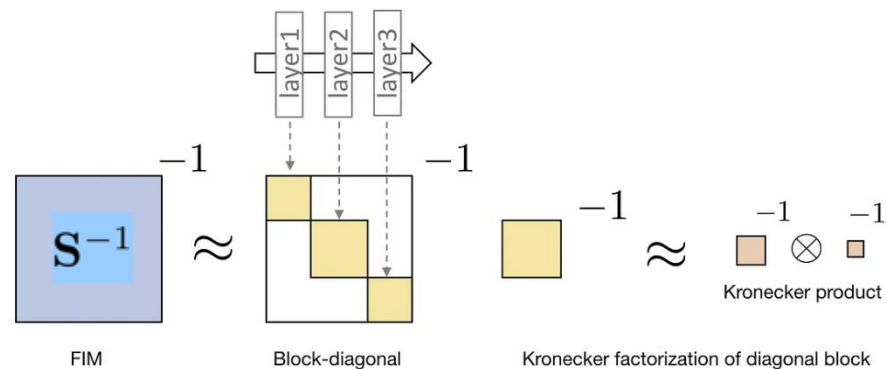
$$S_{m,n} = \left\langle \left(\frac{\partial \ln[\psi(\mathbf{x}; \boldsymbol{\theta})]}{\partial \theta_m} \right)^* \frac{\partial \ln[\psi(\mathbf{x}; \boldsymbol{\theta})]}{\partial \theta_n} \right\rangle - \left\langle \frac{\partial \ln[\psi(\mathbf{x}; \boldsymbol{\theta})]}{\partial \theta_m} \right\rangle^* \left\langle \frac{\partial \ln[\psi(\mathbf{x}; \boldsymbol{\theta})]}{\partial \theta_n} \right\rangle$$

Matrix of size

Number of parameters \times Number of parameters

Neural network quantum states: approaches to train larger networks

(1) K-FAC approximates the inverse by turning it into a block diagonal.

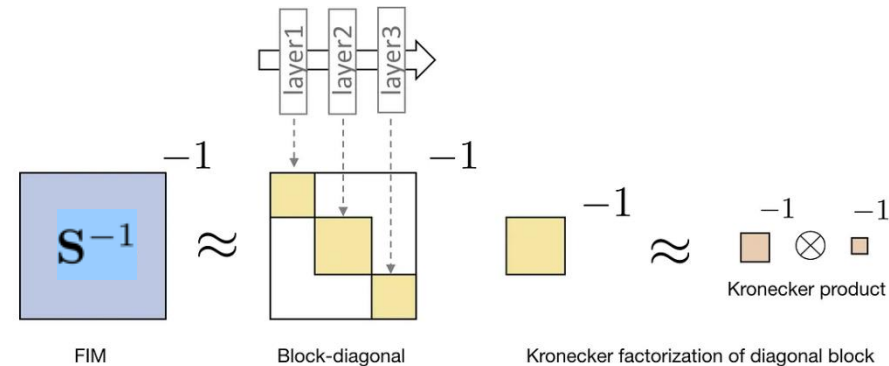


J. Martens and R. Grosse, Proc. Mach. Lear. Res. 37,2408 (2015)

From towardsdatascience.com

Neural network quantum states: approaches to train larger networks

(1) **K-FAC** approximates the inverse by turning it into a block diagonal.



J. Martens and R. Grosse, Proc. Mach. Lear. Res. 37,2408 (2015)

From towardsdatascience.com

(2) **minSR** turns the matrix to invert from being limited by the number of parameters to be limited by the number of samples which can be much smaller than the number of parameters.

$$\begin{aligned}
 (\mathbf{S} + \lambda \mathbb{1}_p)^{-1} \mathbf{F}^U &= (\mathbf{X}\mathbf{X}^\dagger + \lambda \mathbb{1}_p)^{-1} \mathbf{X} \mathbf{f}_E^U \\
 &= \mathbf{X} (\mathbf{X}^\dagger \mathbf{X} + \lambda \mathbb{1}_s)^{-1} \mathbf{f}_E^U
 \end{aligned}$$

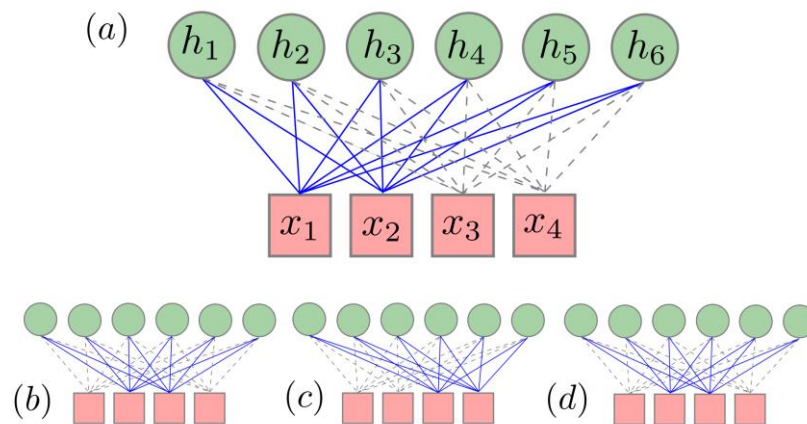
Matrix of size

Number of samples \times Number of samples

Chen and Heyl, Nature Physics 2024
Rende, et al., Comm. Physics 2024

Neural network quantum states: approaches to train larger networks

(3) Sequential Local Optimization SLO is a DMRG-inspired approach which we have developed for ground state search



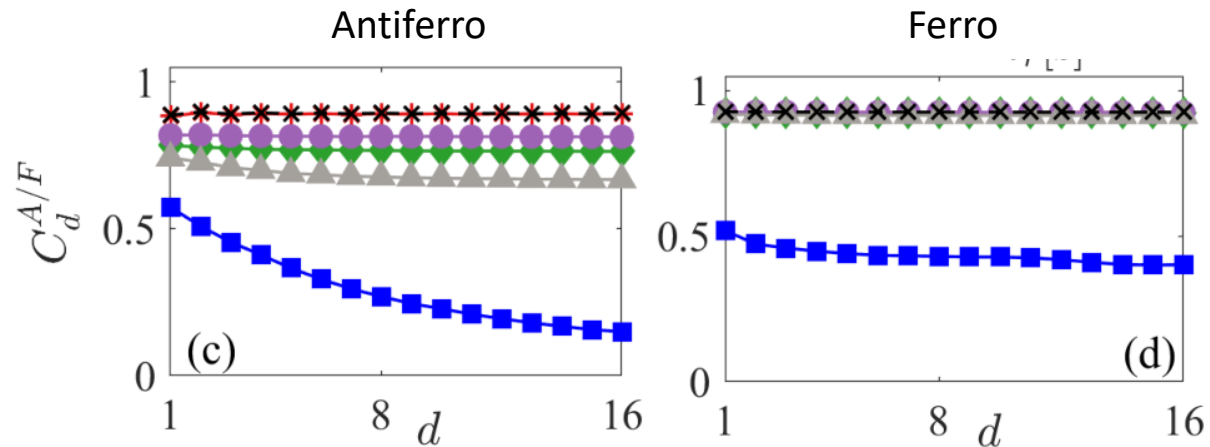
where the key is to optimize only a portion of the network, but to also ensure **overlap between consecutively optimized portions**.

Neural network quantum states: approaches to train larger networks

For example, we consider an Ising model whose ground state can be Ferromagnetic, Antiferromagnetic or Paramagnetic.

$$H = \sum_{\langle i,j \rangle} J \sigma_i^z \sigma_j^z - \sum_i (h_z \sigma_i^z + h_x \sigma_i^x)$$

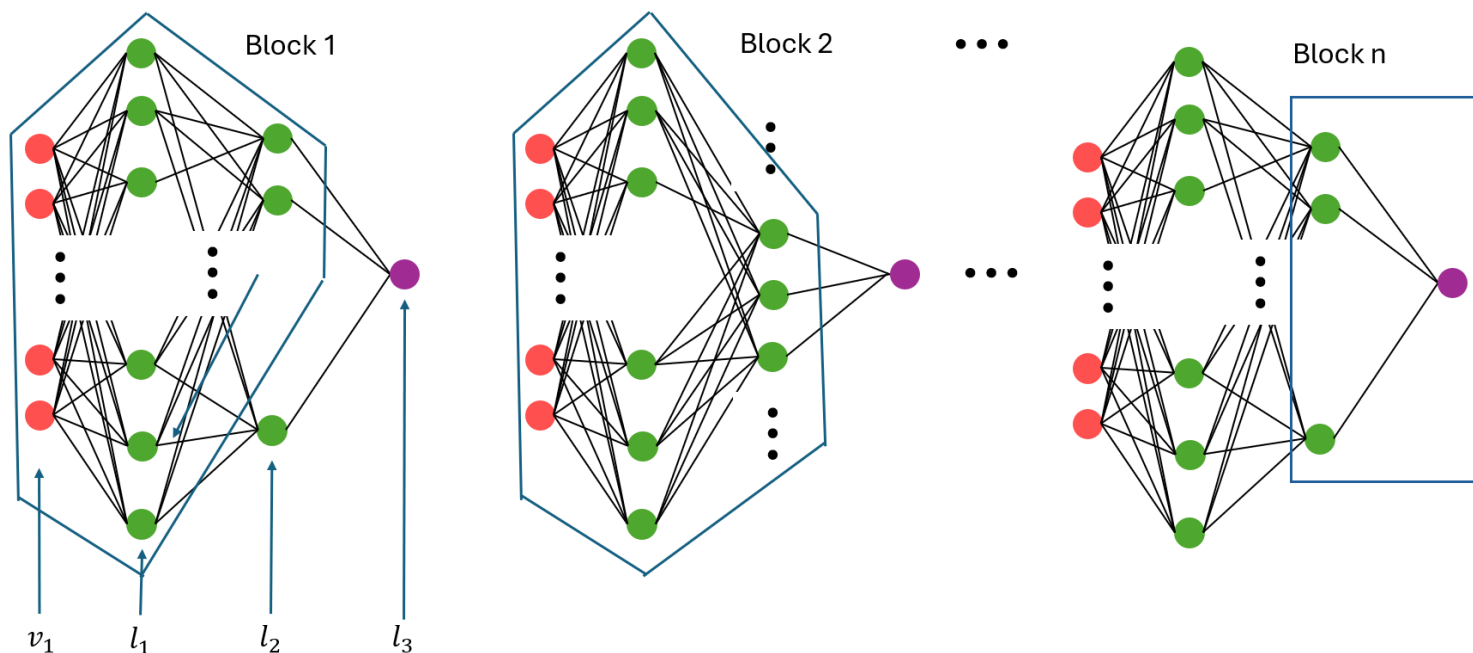
Antiferro and Ferro correlations after 400s



The symbol * represents 2-site block, \blacklozenge 4-site block, \bullet 8-site block, \blacktriangle 16-site block and \blacksquare represents updating all parameters.

Neural network quantum states: approaches to train larger networks

For the time evolution, with our FFNN, we use this type of SLO which resulted to be more stable.



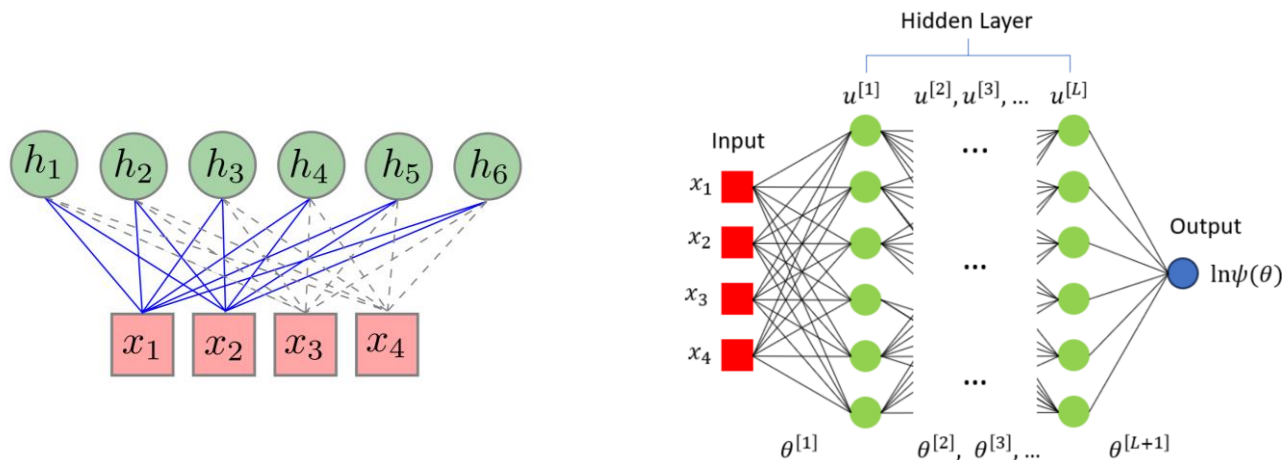
Plan of the presentation

- 1) Time evolution from overlap
 - Loss function
 - Optimization approaches

2) Holomorphic neural networks

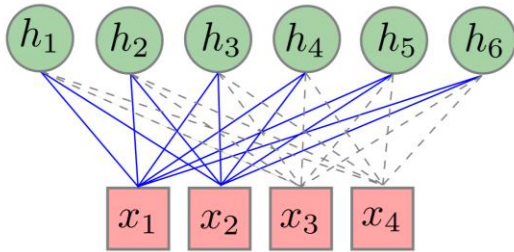
3) Quantum quench in the tilted Ising model

4) Conclusions and outlook



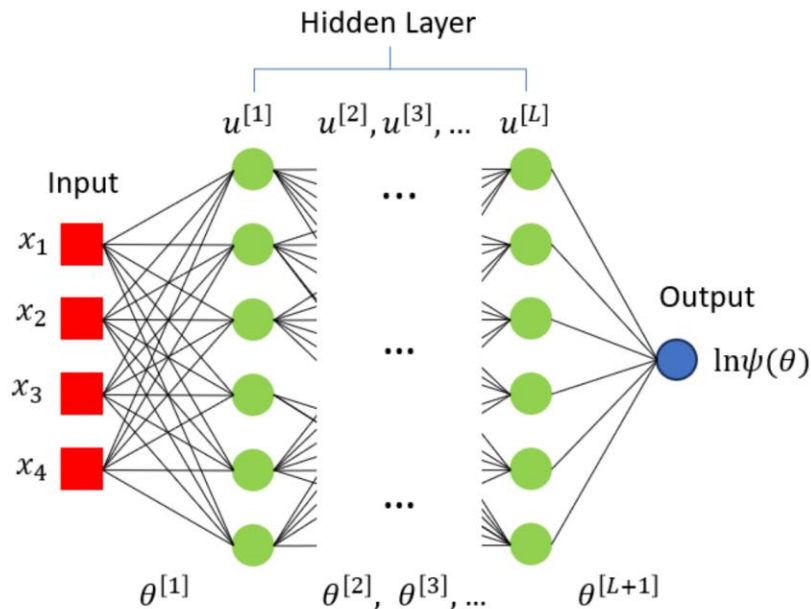
Neural network quantum states considered and their properties

Restricted Boltzmann machines



$$\psi(\mathbf{x}; \boldsymbol{\theta})_{\text{RBM}} = \exp\left(\sum_l a_l x_l\right) \prod_{h=1}^H 2 \cosh\left(b_h + \sum_{l=1}^L W_{h,l} x_l\right)$$

Feed-forward neural networks



$$\mathbf{u}^{[k]} = \mathbf{W}^{[k]} f(\mathbf{u}^{[k-1]}) + \mathbf{b}^{[k]}$$

$$f(u) = u - \frac{1}{3}u^3 + \frac{2}{15}u^5$$

$$\psi(\mathbf{x}; \boldsymbol{\theta})_{\text{FNN}} = \exp\left(\mathbf{u}^{[K]}\right)$$

Neural network quantum states considered and their properties

Our main object of interest is the loss function

$$\mathcal{F}(|\psi_\theta\rangle, |\phi\rangle) = \frac{\langle\psi|\phi\rangle\langle\phi|\psi\rangle}{\langle\psi|\psi\rangle\langle\phi|\phi\rangle}$$

Neural network quantum states considered and their properties

Our main object of interest is the loss function

$$\mathcal{F}(|\psi_\theta\rangle, |\phi\rangle) = \frac{\langle \psi | \phi \rangle \langle \phi | \psi \rangle}{\langle \psi | \psi \rangle \langle \phi | \phi \rangle}$$

Evolved past wave function

Neural network quantum states considered and their properties

Our main object of interest is the loss function

$$\mathcal{F}(|\psi_\theta\rangle, |\phi\rangle) = \frac{\langle \psi | \phi \rangle \langle \phi | \psi \rangle}{\langle \psi | \psi \rangle \langle \phi | \phi \rangle}$$

$$\frac{\partial}{\partial z^*} (f^* f) = \left(\frac{\partial}{\partial z^*} f^* \right) f$$

$$\begin{aligned} \partial_{\theta^*} \mathcal{F} &= \frac{\partial}{\partial \theta^*} \left(\frac{\langle \psi | \phi \rangle \langle \phi | \psi \rangle}{\langle \psi | \psi \rangle \langle \phi | \phi \rangle} \right) \\ &= \frac{\frac{\partial \langle \psi |}{\partial \theta^*} | \phi \rangle}{\langle \psi | \psi \rangle} \cdot \frac{\langle \phi | \psi \rangle}{\langle \phi | \phi \rangle} - \frac{\langle \psi | \phi \rangle \langle \phi | \psi \rangle}{\langle \phi | \phi \rangle} \cdot \frac{1}{(\langle \psi | \psi \rangle)^2} \cdot \frac{\partial \langle \psi |}{\partial \theta^*} | \psi \rangle \end{aligned}$$



$$\partial_{\theta^*} \mathcal{F} = \mathbb{E}_{\mathbf{y} \sim \pi_\phi} \left[\frac{\psi(\mathbf{y})}{\phi(\mathbf{y})} \right] \cdot \mathbb{E}_{\mathbf{x} \sim \pi_\psi} \left[(\Delta J(\mathbf{x}))^* \cdot \frac{\phi(\mathbf{x})}{\psi(\mathbf{x})} \right]$$

$$\Delta J(\mathbf{x}) = J(\mathbf{x}) - \mathbb{E}_{\mathbf{z} \sim \pi_\psi} [J(\mathbf{z})]$$

$$J(\mathbf{x}) = \partial_{\theta} \log \psi(\mathbf{x})$$

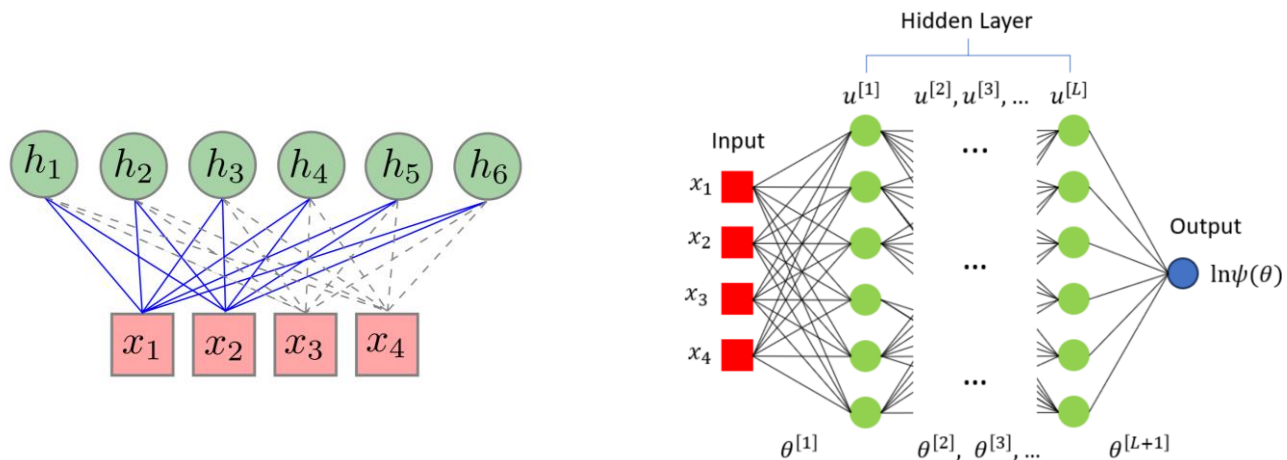
Plan of the presentation

- 1) Time evolution from overlap
 - Loss function
 - Optimization approaches

2) Holomorphic neural networks

3) Quantum quench in the tilted Ising model

4) Conclusions and outlook



Quantum quench in the tilted Ising model

Tilted Ising model (TIM)

$$H = J \sum_{\langle i,j \rangle} \sigma_i^z \sigma_j^z - h_x \sum_k \sigma_k^x - h_z \sum_m \sigma_m^z$$

Hamiltonian is divided into blocks:

$$H = \sum_p H_p$$

with

$$H_p = J \sigma_p^z \sigma_q^z - h_x (n_p \sigma_p^x + n_q \sigma_q^x) - h_z (n_p \sigma_p^z + n_q \sigma_q^z)$$

n_α with $\alpha = p, q$ is coordination number of the site and for 1D TIM

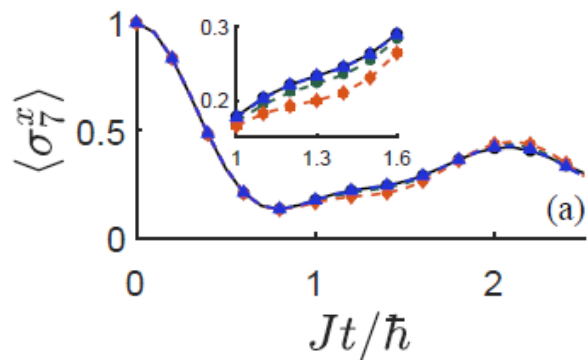
$$\begin{cases} n_\alpha = 1, \alpha \in \{1, L\} \\ n_\alpha = \frac{1}{2}, \alpha \notin \{1, L\} \end{cases}$$

The unitary operator can include more than one local operator

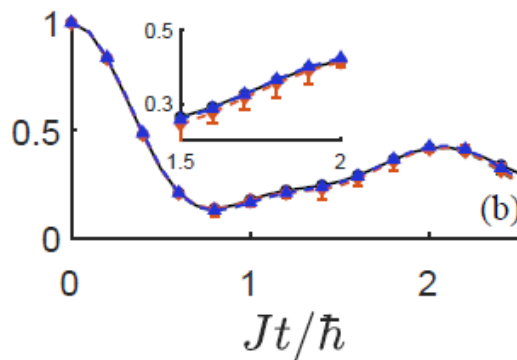
$$U = \exp(-i \sum_{\{p\}} H_p \cdot dt/2)$$

Quantum quench in the tilted Ising model

FFNN [14,56,42,1]



RBM $\alpha=5$



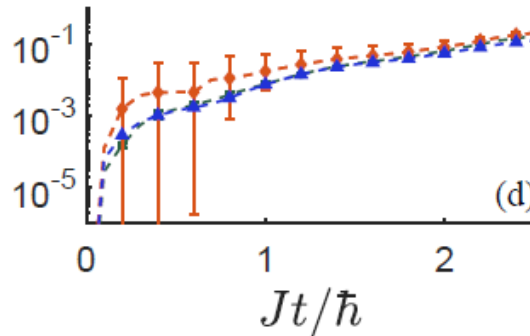
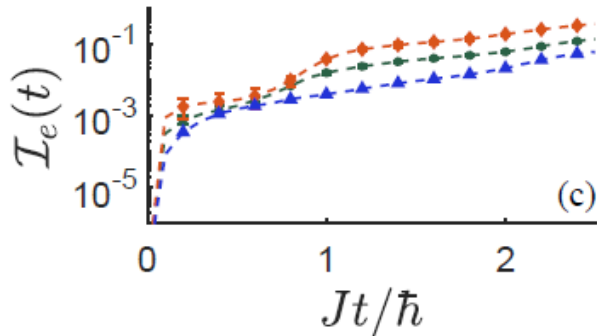
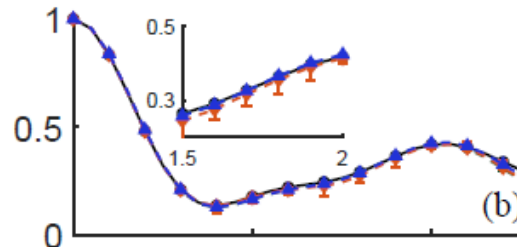
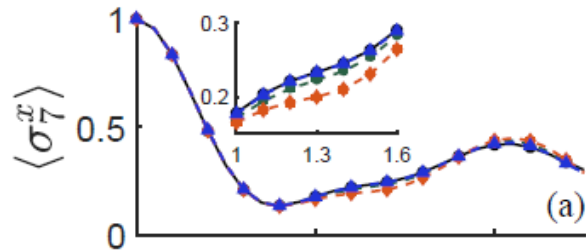
L = 14

- Exact Diagonalization
- ◇ Runge-Kutta ($dt=10^{-3}$)
- * Overlap ($dt=10^{-1}$)
- △ Overlap Full summation ($dt=10^{-1}$)

Quantum quench in the tilted Ising model

FFNN [14,56,42,1]

RBM $\alpha=5$



Cumulated error from exact state

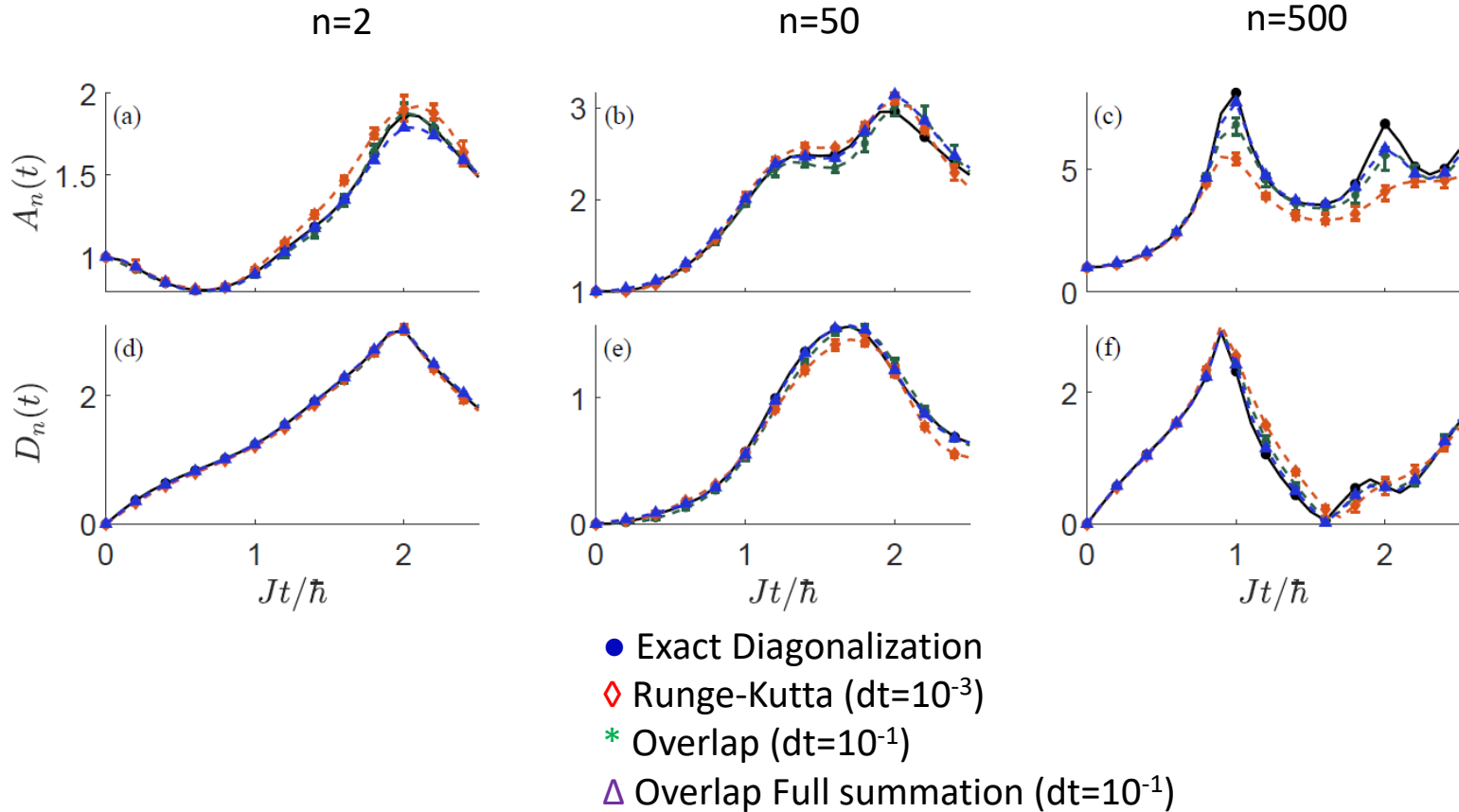
$$I_e(t) = \left| 1 - \frac{\langle \psi_t | \psi_{e,t} \rangle \langle \psi_{e,t} | \psi_t \rangle}{\langle \psi_t | \psi_t \rangle \langle \psi_{e,t} | \psi_{e,t} \rangle} \right|$$

$$\mathcal{I}_e(t) = \int_0^t I_e(t) dt$$

$L = 14$

- Exact Diagonalization
- ◇ Runge-Kutta ($dt=10^{-3}$)
- * Overlap ($dt=10^{-1}$)
- △ Overlap Full summation ($dt=10^{-1}$)

Quantum quench in the tilted Ising model



Ratio of amplitude between most likely and n-th likely configuration

$$A_n(t) = \left| \frac{\psi_t(\mathbf{x}_M)}{\psi_t(\mathbf{y}_n)} \right|$$

Relative phase between most likely and n-th likely configuration

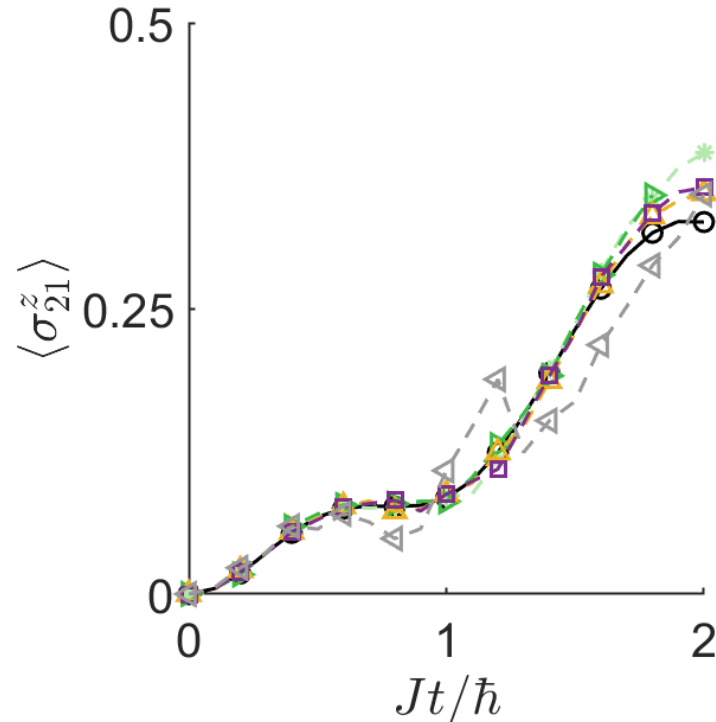
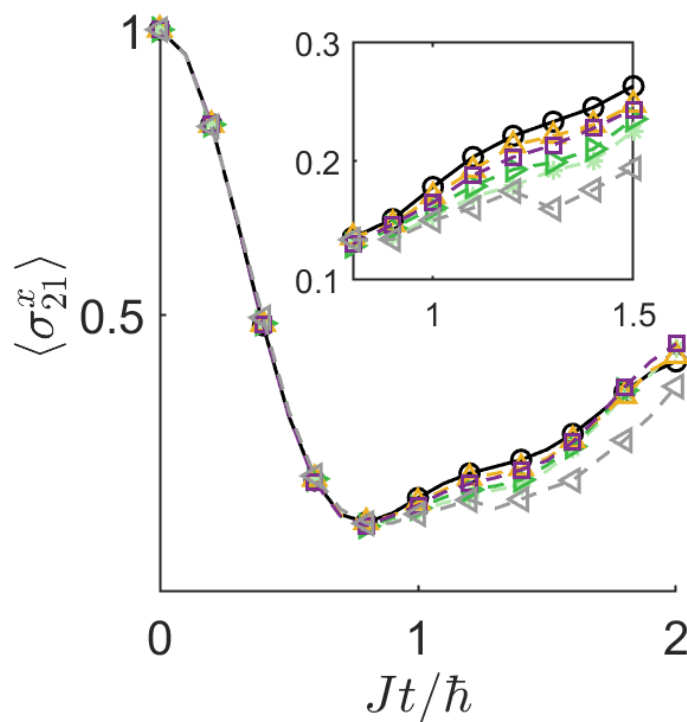
$$D_n(t) = \min(d_n(t), 2\pi - d_n(t))$$

$$d_n(t) = |\text{Arg}(\psi_t(\mathbf{x}_M)) - \text{Arg}(\psi_t(\mathbf{y}_n))|$$

Quantum quench in the tilted Ising model

$L = 40$

FFNN [L, 4L, 3L, 1]



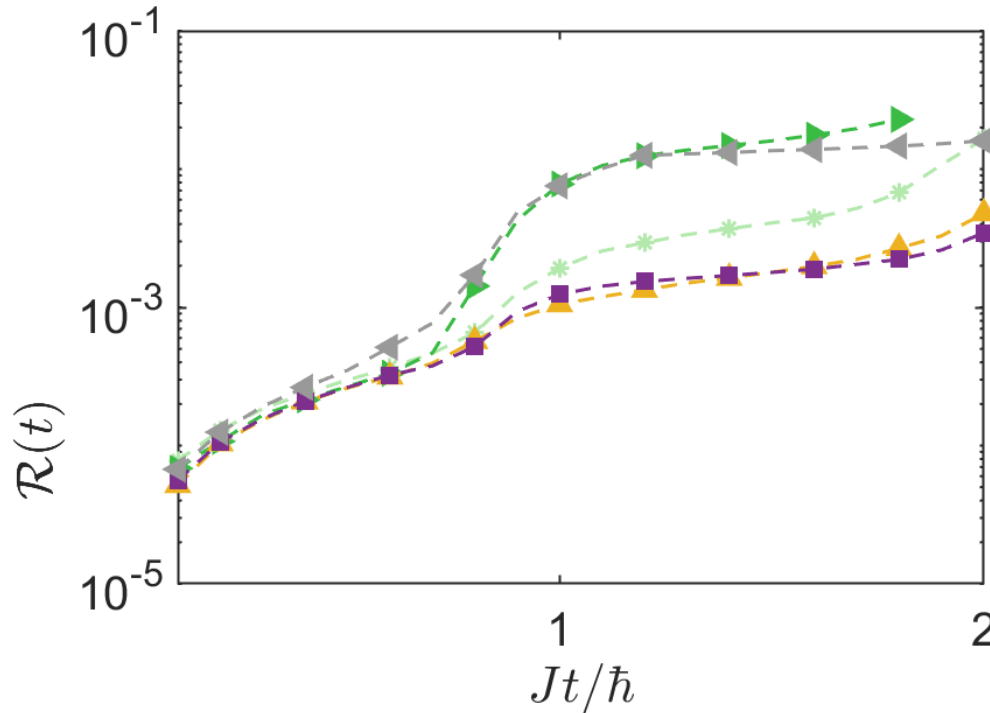
- t-MPS
- * K-FAC with 10^4 samples
- ▷ K-FAC with 10^5 samples
- SLO with 10^4 samples
- △ minSR with 10^4 samples
- ◁ AdamW with 10^4 samples

- minSR seems to perform better
- SLO is more accurate than K-FAC
- K-FAC does not seem to improve significantly with more sample

Quantum quench in the tilted Ising model

L = 40

FFNN [L, 4L, 3L, 1]



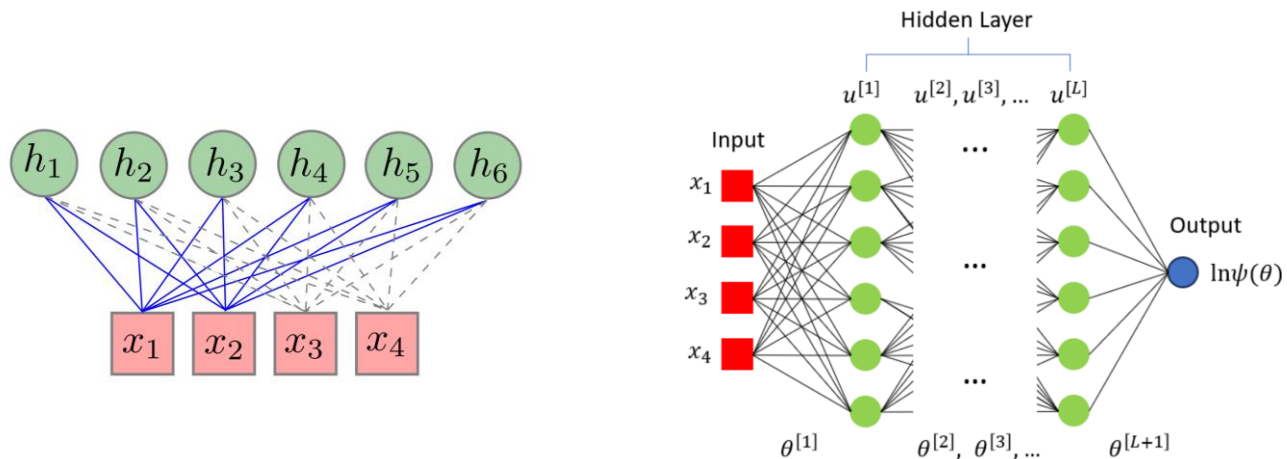
Cumulated error from state we have

$$\mathcal{R}(t) = \int_0^t \sum_l I_{\psi_t, \psi'_t}^{U_{d,l}} dt$$

- t-MPS
 - * K-FAC with 10^5 samples
 - ▷ K-FAC with 10^4 samples
 - ◻ SLO with 10^4 samples
 - △ minSR with 10^4 samples
 - ◁ AdamW with 10^4 samples
- minSR seems to perform better
 - SLO is more accurate than K-FAC
 - K-FAC does not seem to improve significantly with more sample

Plan of the presentation

- 1) Time evolution from overlap
 - Loss function
 - Optimization approaches
- 2) Holomorphic neural networks
- 3) Quantum quench in the tilted Ising model
- 4) Conclusions and outlook



Conclusions and outlook

Ground state search-like methods can be used for time evolution, like minSR, K-FAC and SLO.

minSR seems, in our limited exploration, the better performing and more stable one.

Comparable performance can be achieved by SLO.



Conclusions and outlook

Ground state search-like methods can be used for time evolution, like minSR, K-FAC and SLO.

minSR seems, in our limited exploration, the better performing and more stable one.

Comparable performance can be achieved by SLO.

Outlook:

- Translate this to “lighter” and more expressive networks
- Further understand the role of sampling and how control variate can help (see latest work by Gravina, Savona and Vicentini)



Thank you!

Principal Investigator
Dario Poletti

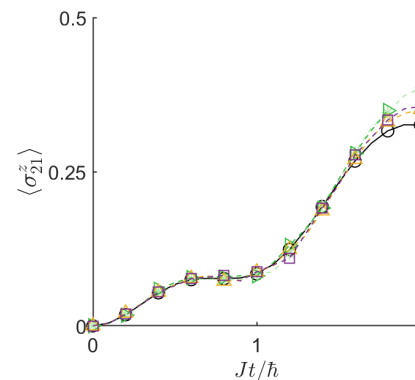
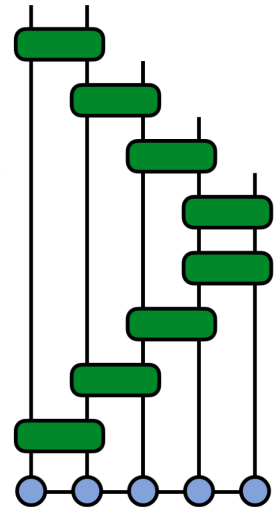
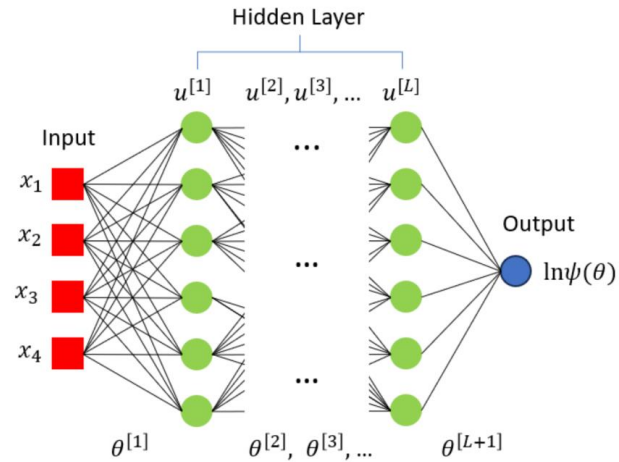
Collaborators on
these topics

PostDocs
Noufal Jaseem
Wang Dingzu
Zhang Wenxuan

Guo Chu
Itai Arad
Lu Wei
Kavan Modi

PhD Students
Rebecca Erbanni
Gauthameswar
Apimuk Sornsang

Remmy Zen
Stephane Bressan
Christian Miniatura
Xu Xiansong
Xing Bo



Zhang et al., arxiv:2406.03381